

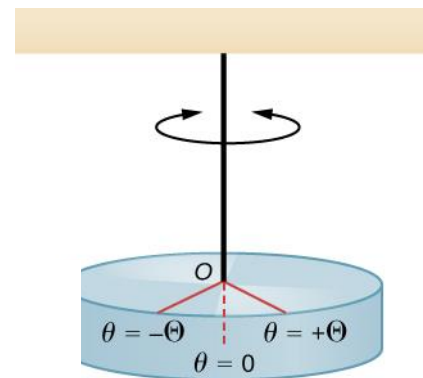
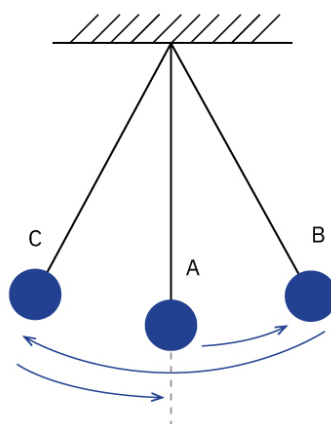
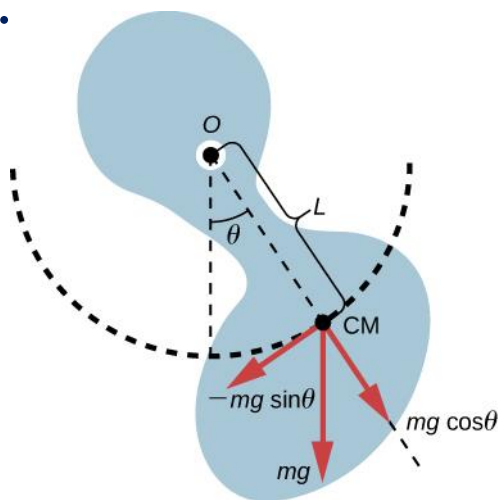
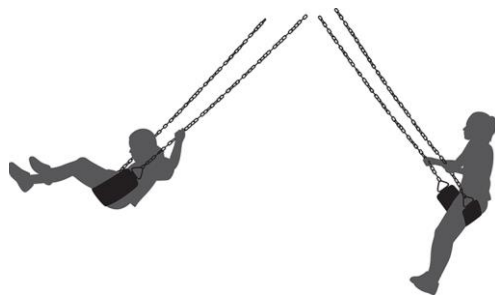


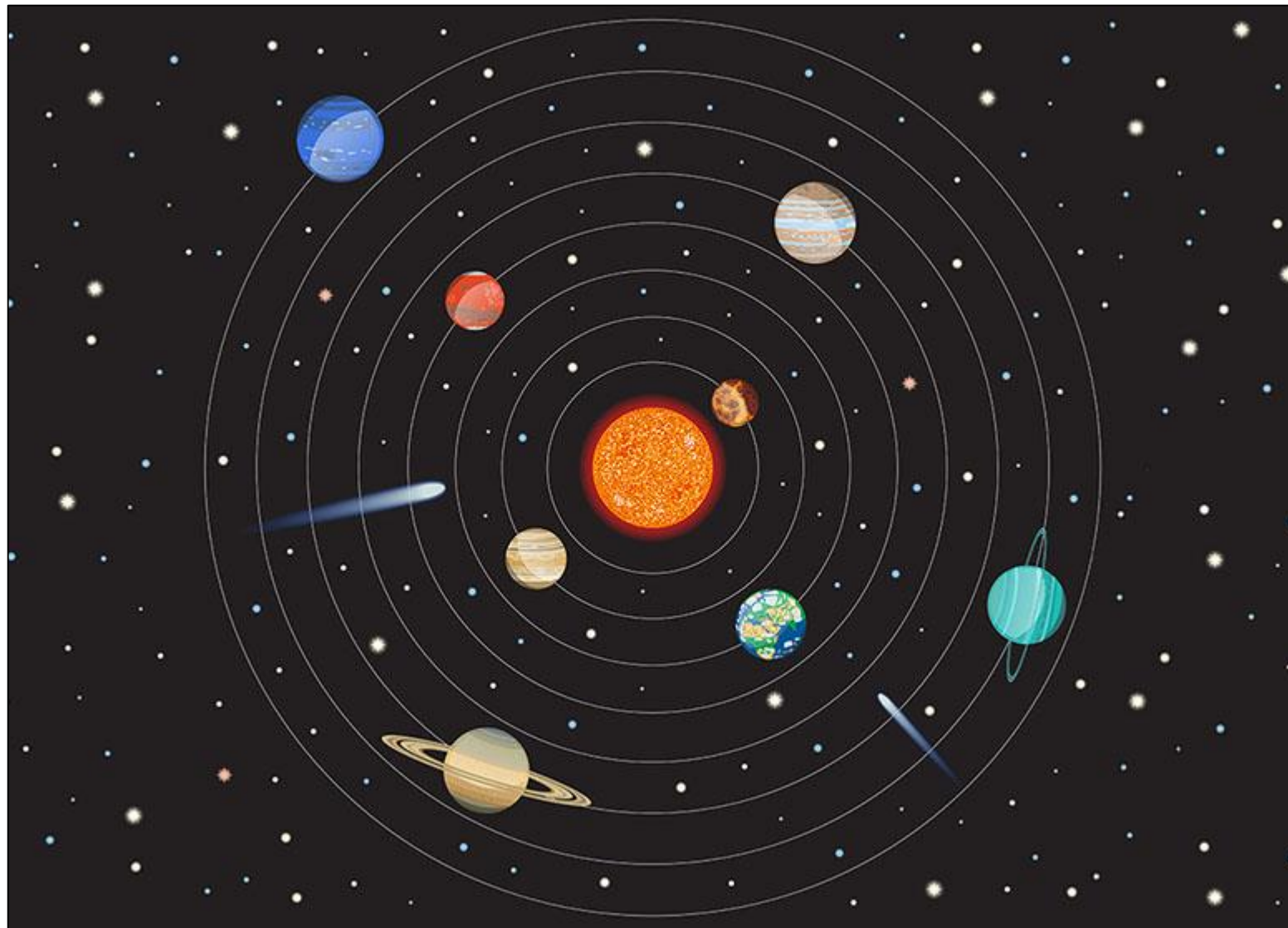
Fizika I

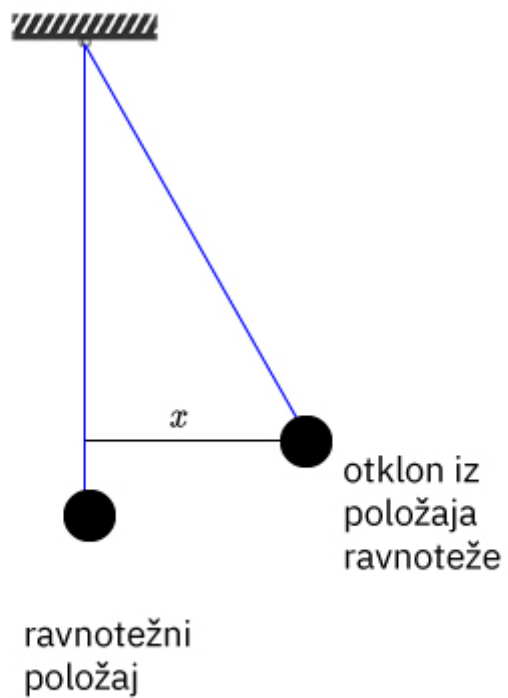
TITRANJE

TITRANJE

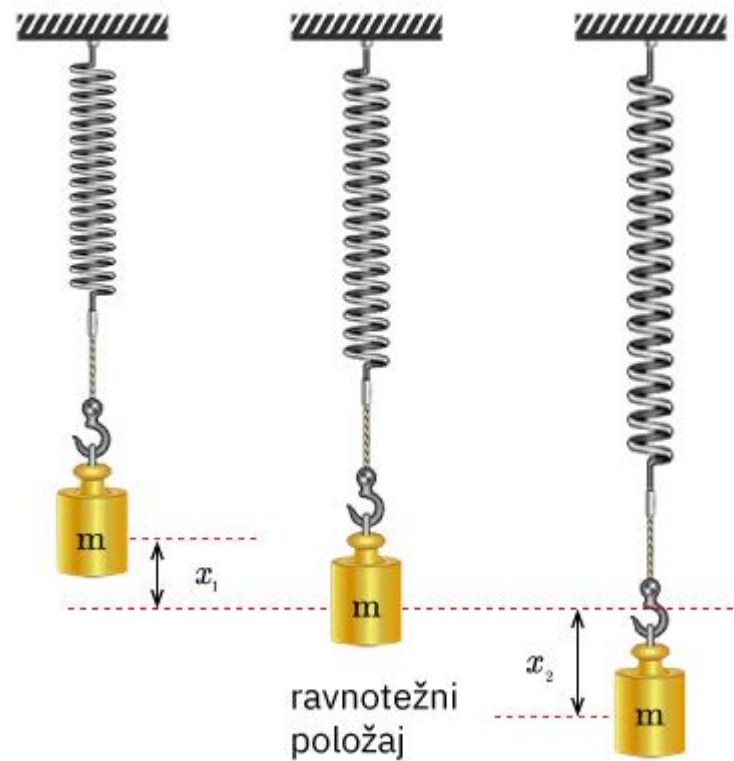
- ▶ Titranje (**oscilatorno gibanje**) - gibanje pri kojem materijalna čestica prevaljuje određeni put između dvaju krajnjih položaja, vraća se istim putem nakon čega se to gibanje periodički ponavlja.
- ▶ Čestica koja titra - **oscilator**.



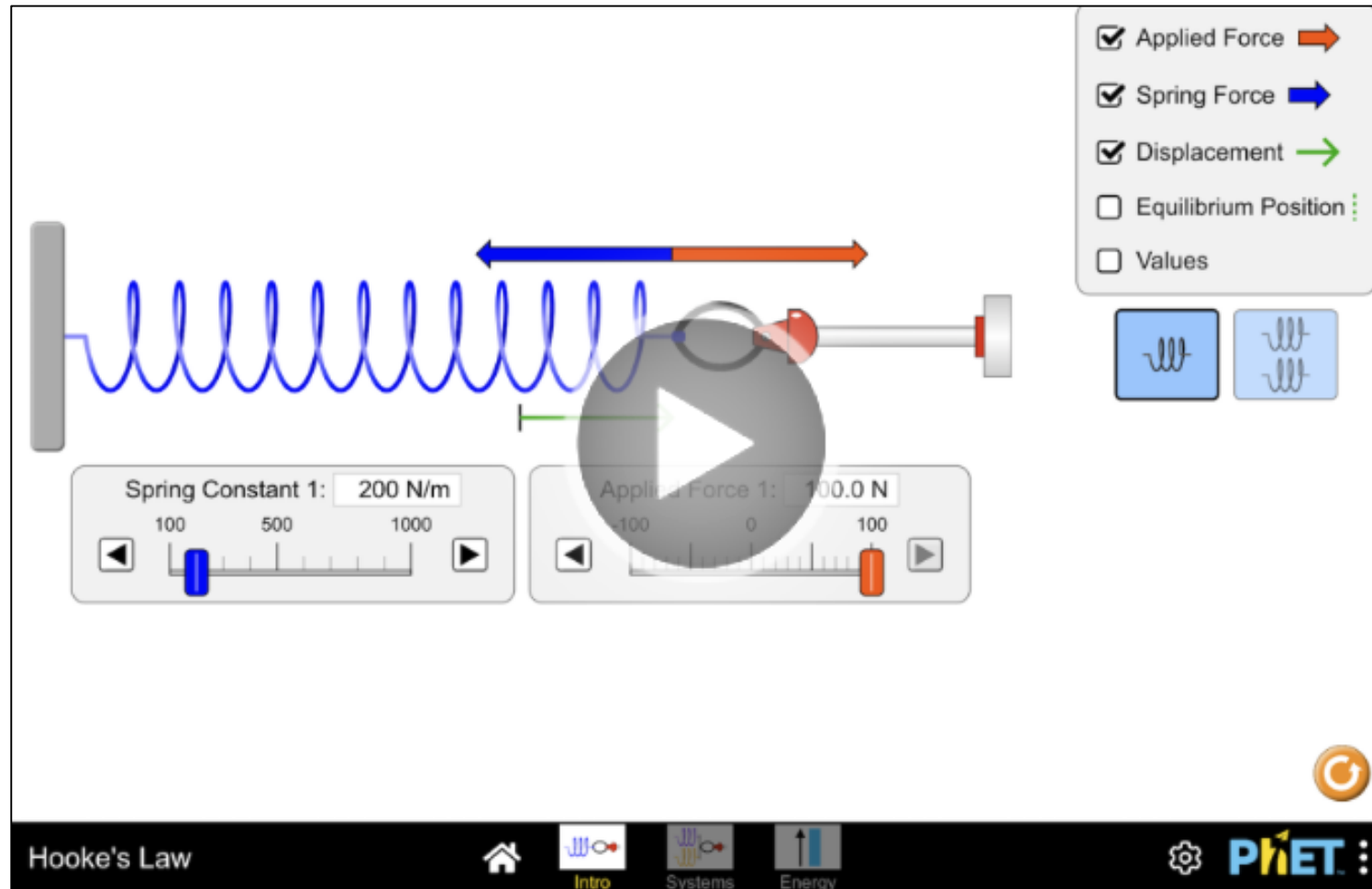




x_1, x_2 - elongacija, odn. pomak iz ravnotežnog položaja



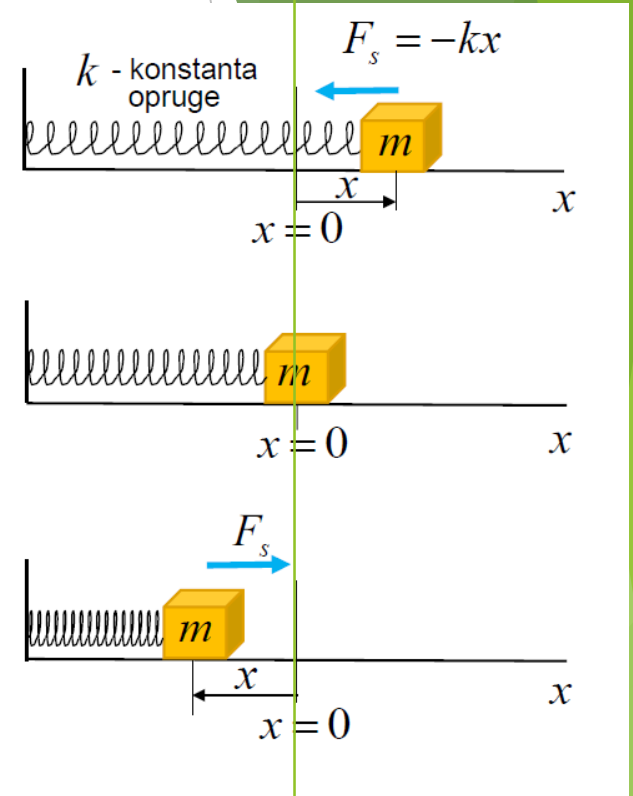
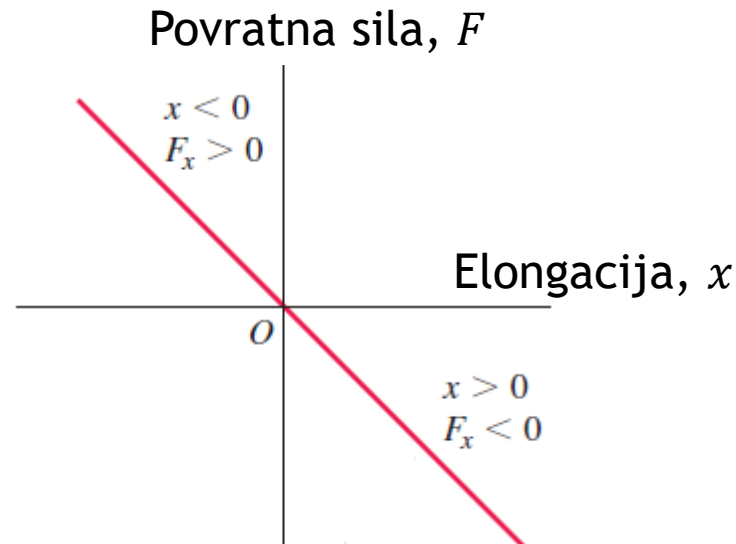
Simulacija: Hookeov zakon



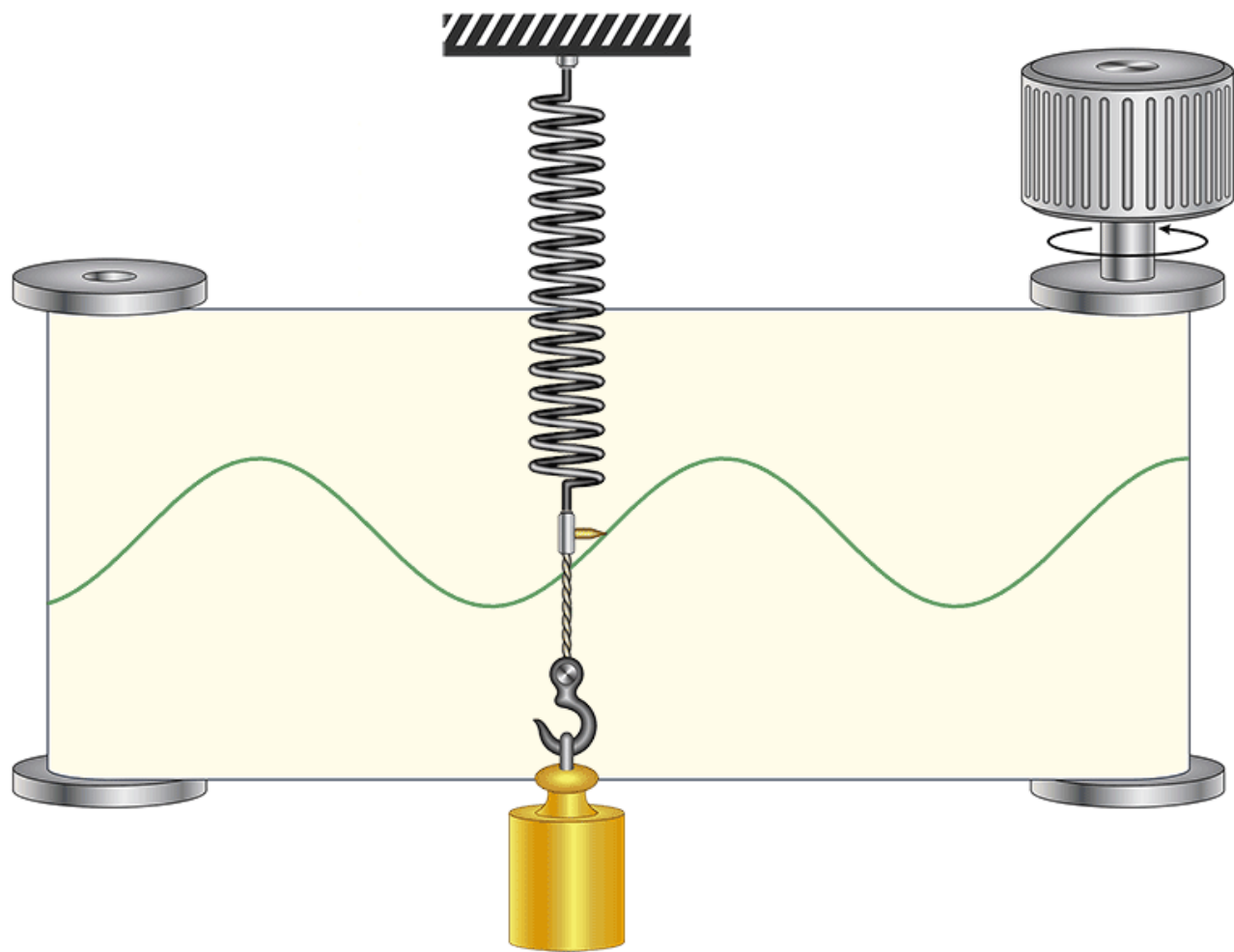
Harmoniĉko titranje

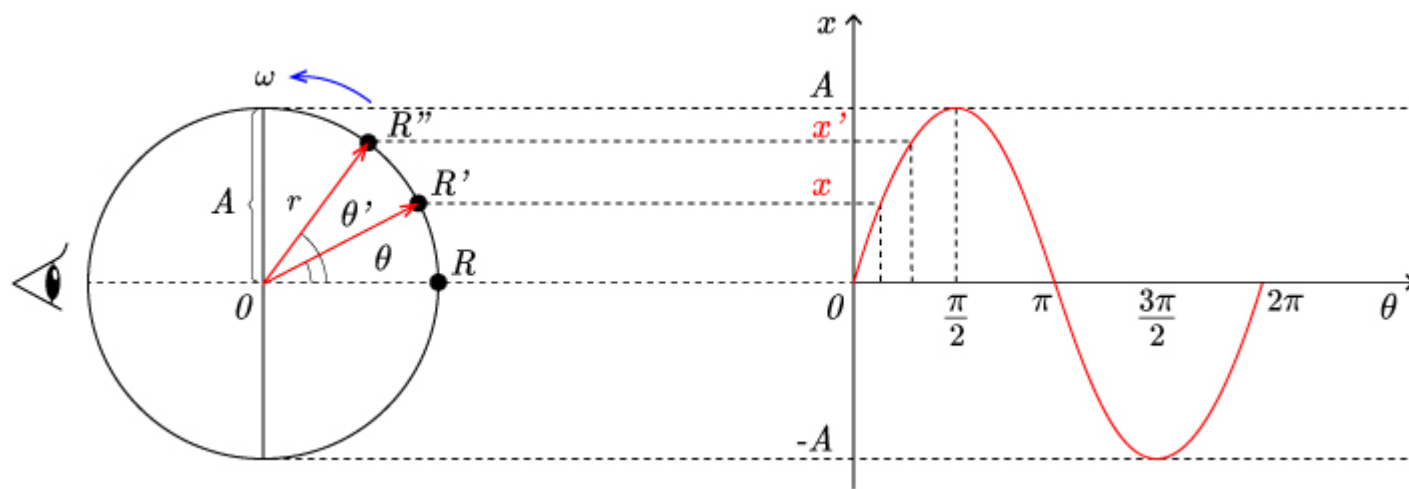
- Posebno pravilan naĉin titranja.
- Ćestica koja harmoniĉki titra - **harmonički oscilator**.
- Uzrok gibanju je elastiĉna (harmonička) sila.
- Elastiĉna sila je **povratna** sila - uvijek suprotna od pomaka ĉestice iz ravnoteŹnog poloŹaja.
- Povratna sila je direktno proporcionalna elongaciji iz ravnoteŹe x (s): (Hookeov zakon).

$$F = -kx$$

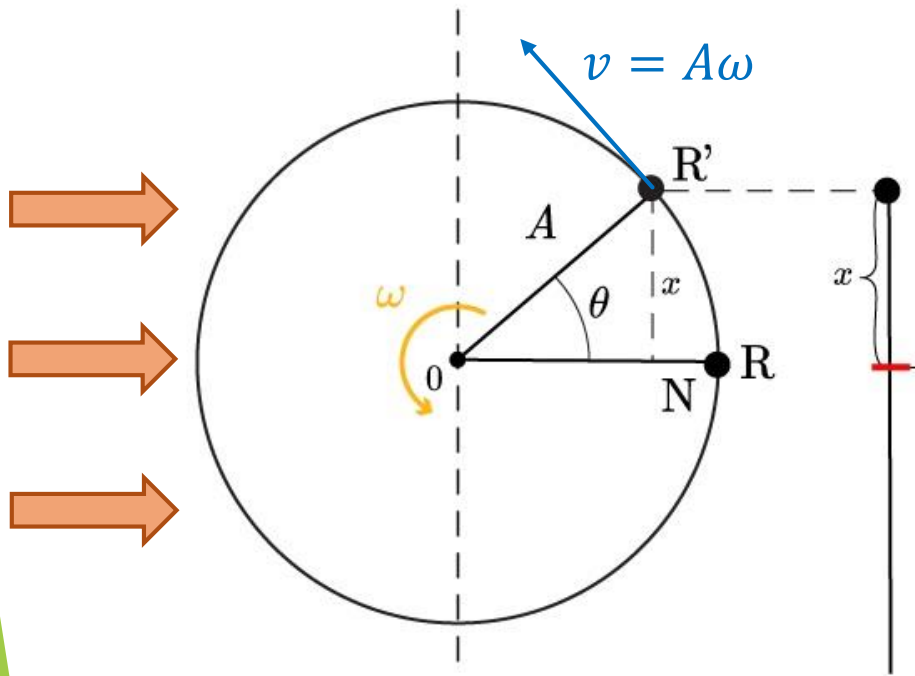


ravnoteŹni
poloŹaj





Harmoniĉko titranje kao projekcija jednolikog gibanja po kruŹnici



- Čestica mase m giba se jednoliko kutnom brzinom $\omega = \theta/t$ u vertikalnoj ravnini po kruŹnici polumjera A .
- Projekcija čestice (**oscilator**) titra gore - dolje

$$x = A \sin \omega t$$

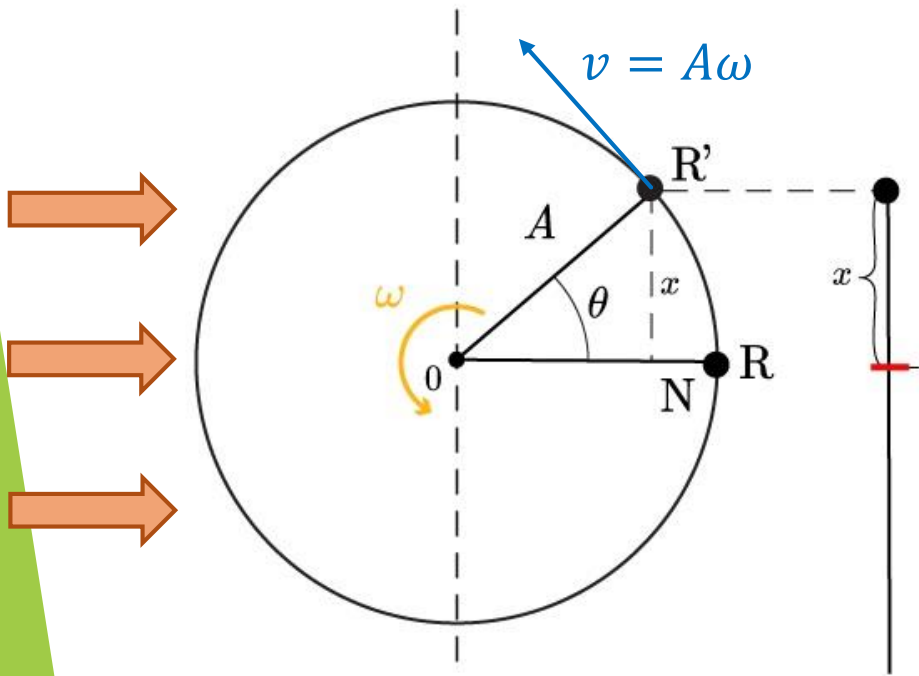
$$v = A \omega \cos \omega t$$

$$a = -\frac{v^2}{A} \sin \omega t$$

$$F = -m \frac{v^2}{A} \sin \omega t = -m \omega^2 x$$

- Povratna harmoniĉka sila: $F = -kx$.

Harmoniĉko titranje kao projekcija jednolikog gibanja po kruŹnici



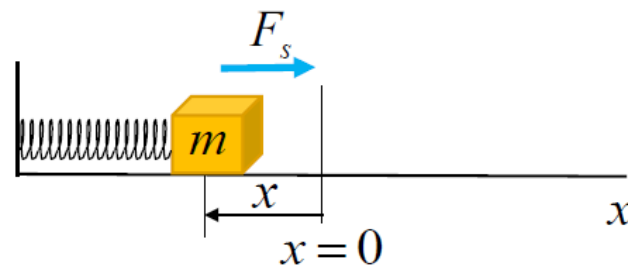
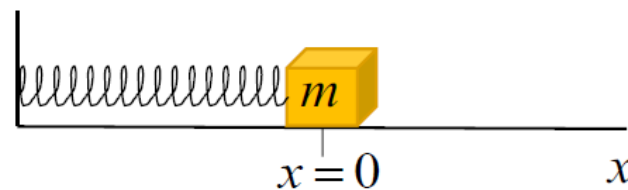
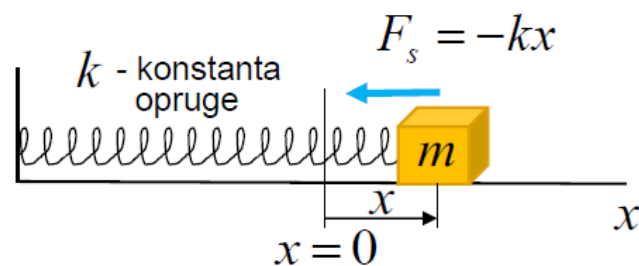
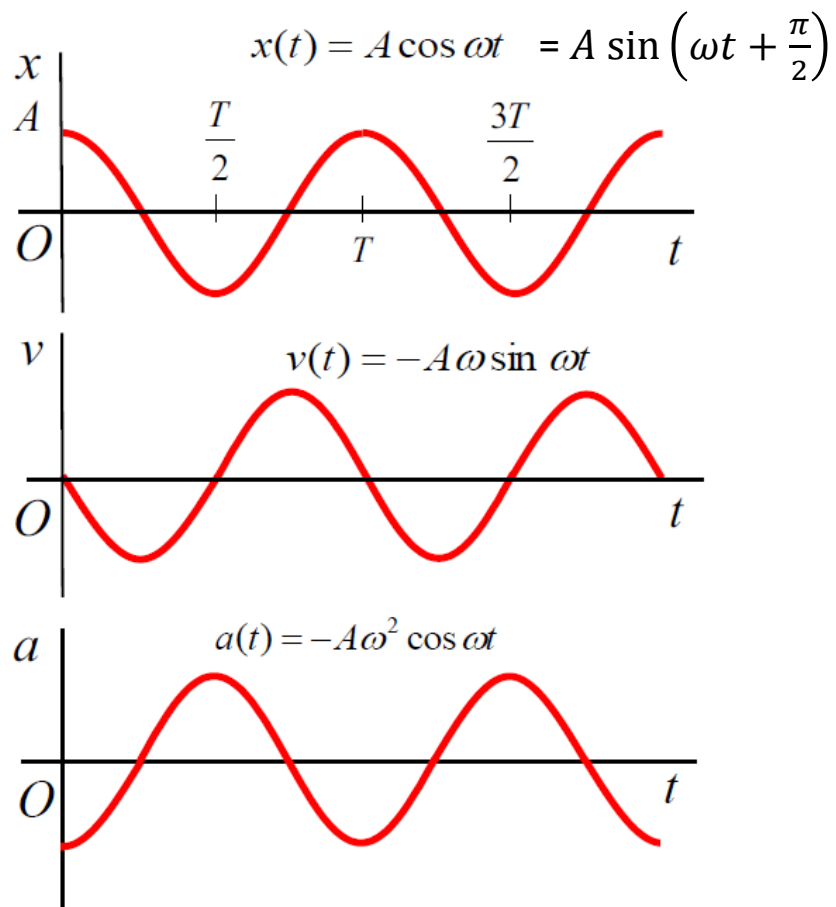
$$k = m\omega^2$$

- KruŹna frekvencija: $\omega = \sqrt{\frac{k}{m}} = 2\pi f$
- Period: $T = 2\pi \sqrt{\frac{m}{k}}$
- Frekvencija: $\nu = f = \frac{1}{T} = \frac{1}{2\pi} \sqrt{\frac{k}{m}}$

$$x = s = A \sin(\omega t + \varphi_0)$$

elongacija početna faza
amplituda

Jednadžba gibanja harmoničkog oscilatora



Hookeov zakon:

$$F = -ks$$

Drugi Newtonov zakon:

$$ma = -ks$$

$$m \frac{d^2 s}{dt^2} = -ks$$

Diferencijalna
jednadžba:

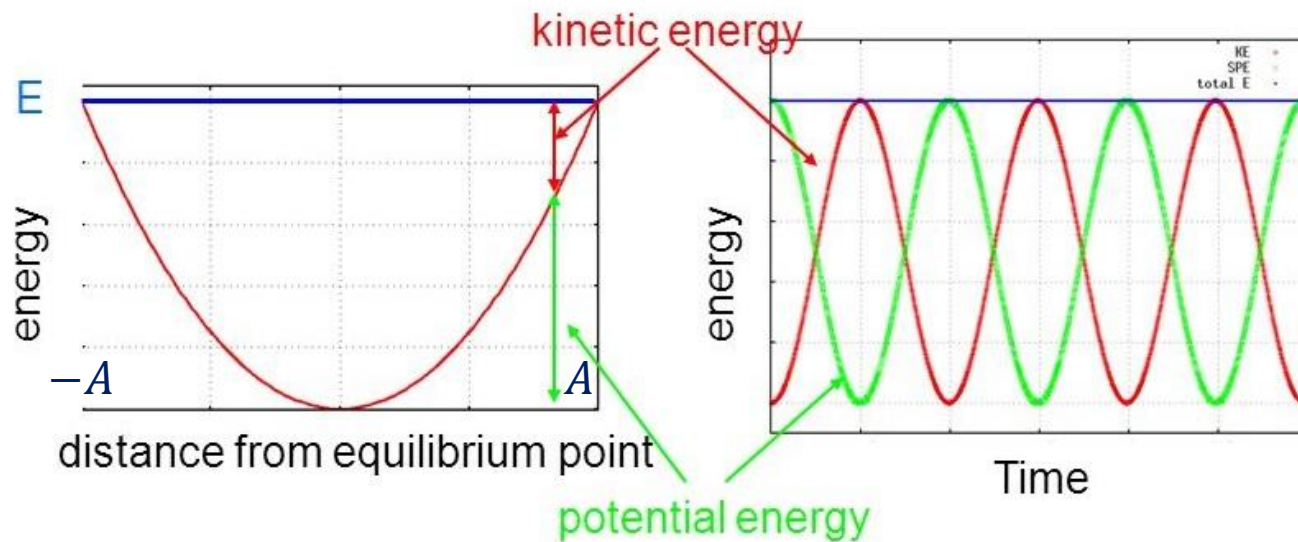
$$\frac{d^2 s}{dt^2} + \underbrace{\frac{k}{m}}_{\omega^2} s = 0$$

Rješenje:

$$s = A \sin(\omega t + \varphi_0)$$

Energija harmoničkog oscilatora

► Simulacija (Geogebra)



Potencijalna energija:

$$E_p = \frac{1}{2} kx^2$$

Kinetička energija:

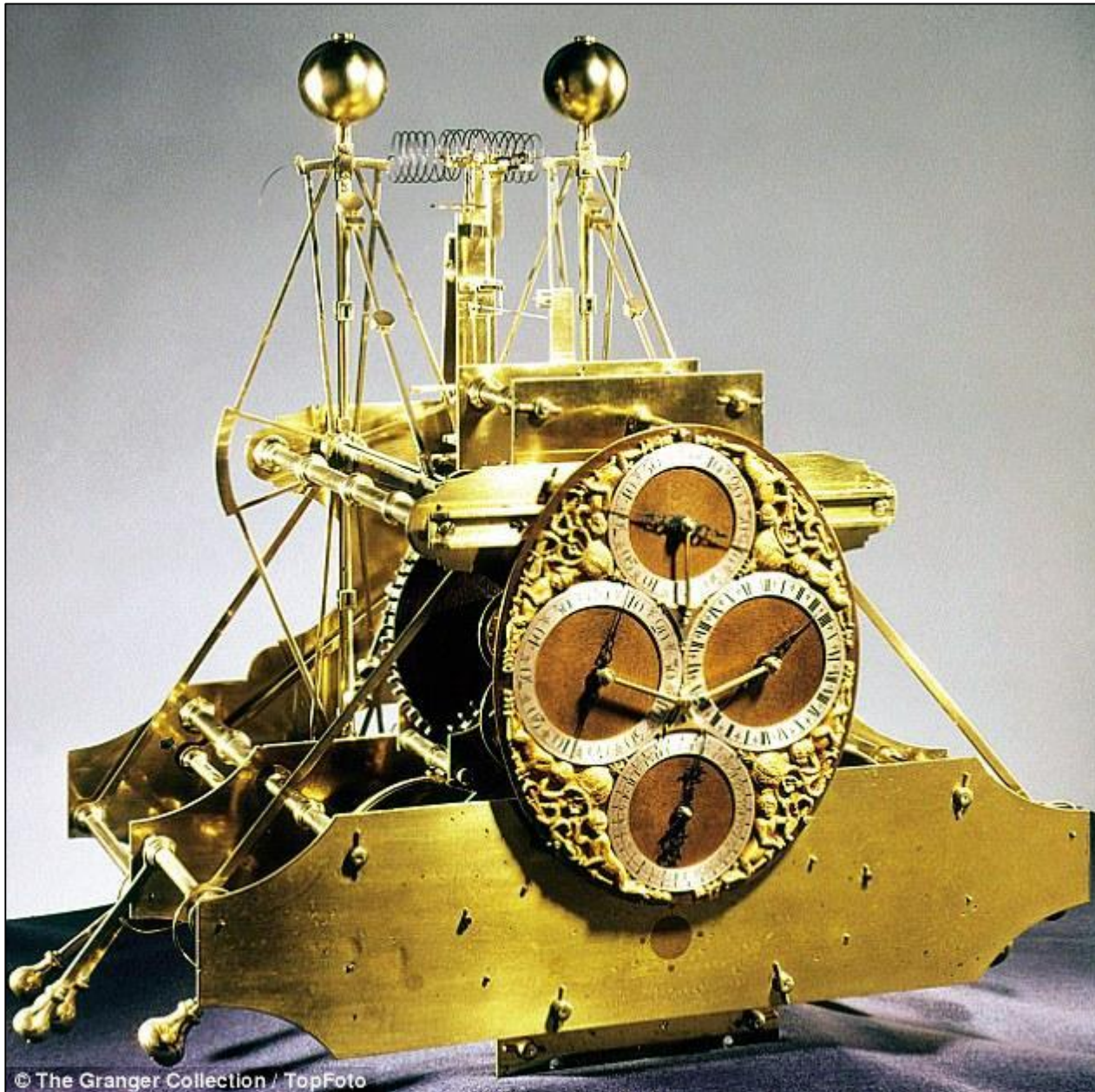
$$E_k = \frac{1}{2} mv^2$$

Ukupna energija

$$E = E_k + E_p$$

$$E = \frac{1}{2} kA^2$$

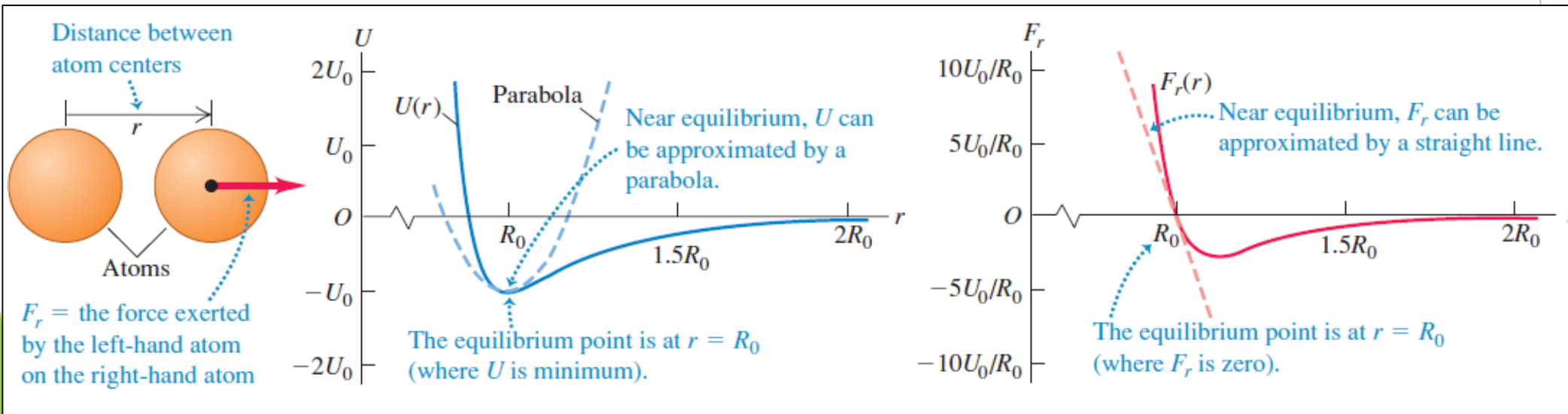




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Vibracije dvoatomne molekule - primjer

- U stvarnosti su mnoga gibanja samo približno harmonička.
- Dvoatomna molekula - anharmonička sila.

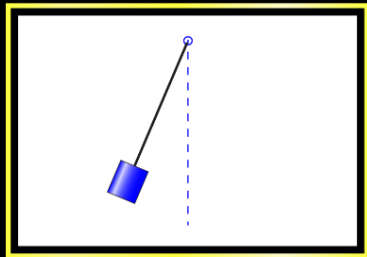


$$U = U_0 \left[\left(\frac{R_0}{r} \right)^{12} - 2 \left(\frac{R_0}{r} \right)^6 \right]$$

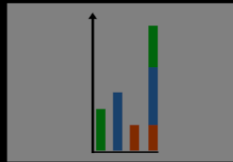
van der Waalsova interakcija

Simulacija: Matematičko njihalo

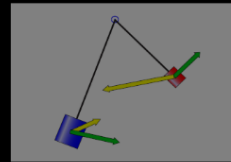
Pendulum Lab



Intro

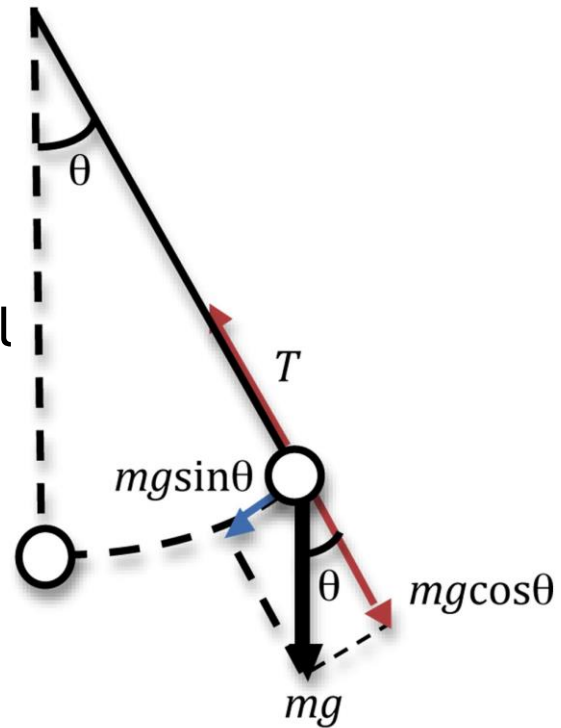


Energy



Lab

Matematičko njihalo



- Idealan sustav: materijalna točka mase m obješena o nerastezljivu nit zanemarive mase i duljine l .
- Dobra aproksimacija: kuglica obješena na kraj duge, tanke nerastezljive niti.
- Kuglica se zbog napete niti giba po kružnici polumjera l .

Newtonov zakon:

$$M = I\alpha = I \frac{d^2\theta}{dt^2}$$

$$I = ml^2$$

$$F_t = mg \sin \theta$$

$$M = -mgl \sin \theta$$

$$\frac{d^2\theta}{dt^2} + \frac{g}{l} \sin \theta = 0$$

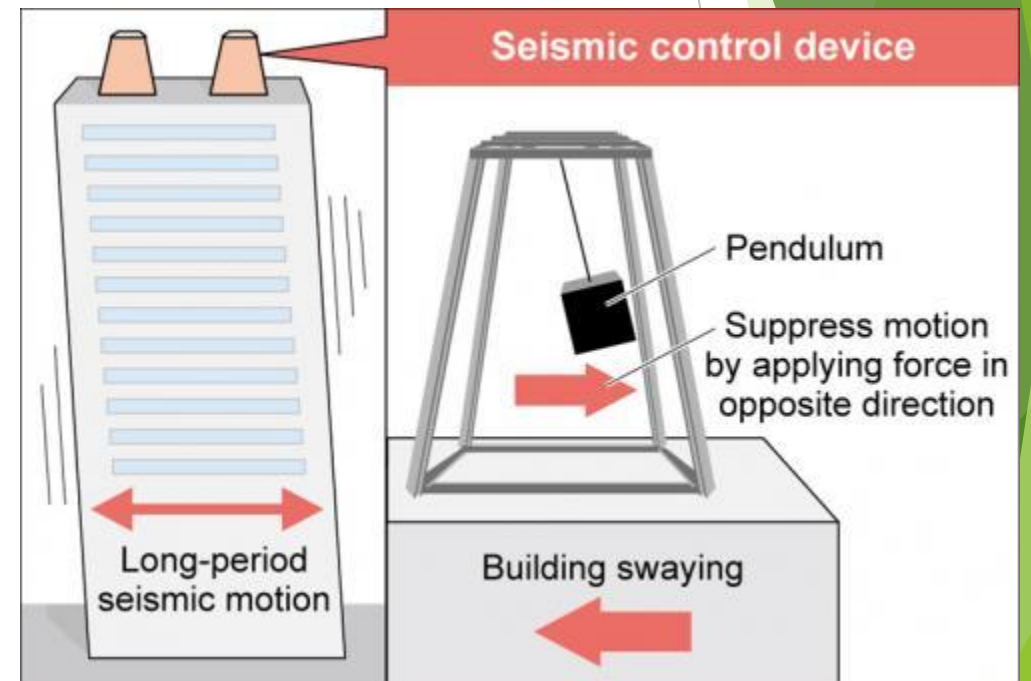
$$\theta \ll 1 \rightarrow \sin \theta \cong \theta$$

Mali kut u radijanima:

$$\frac{d^2\theta}{dt^2} + \left(\frac{g}{l}\right)\theta = 0$$

$$T = 2\pi \sqrt{\frac{l}{g}}$$

$$\theta = \theta_0 \sin(\omega t + \varphi_0)$$



Fizičko njihalo

- ▶ Kruto tijelo obješeno oko čvrste osi koja ne prolazi njegovim težištem.
- ▶ Spojnica između objesišta i težišta (cm) (duljine L) zatvara s težinom $\vec{G}$ kut θ .

Newtonov zakon:

$$M = I\alpha = I \frac{d^2\theta}{dt^2}$$

$$M = -Lmg \sin \theta$$

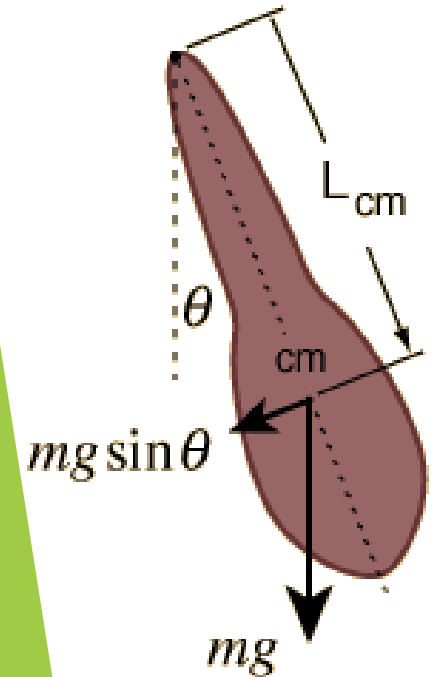
$$\frac{d^2\theta}{dt^2} + \frac{Lmg}{I} \sin \theta = 0$$

$$\theta \ll 1 \rightarrow \sin \theta \cong \theta$$

$$\frac{d^2\theta}{dt^2} + \frac{Lmg}{I} \theta = 0$$

$$T = 2\pi \sqrt{\frac{I}{Lmg}} \quad \leftarrow \omega^2$$

$$\theta = \theta_0 \sin(\omega t + \varphi_0)$$





Prigušeno titranje

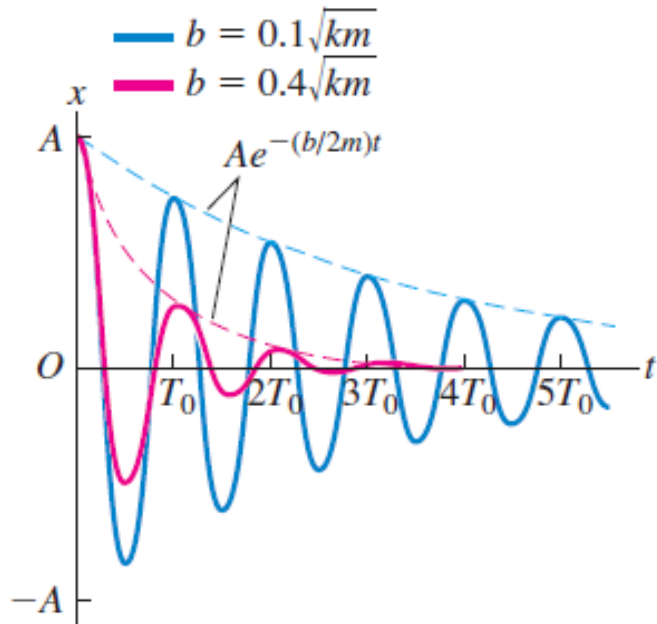
U stvarnim sustavima imamo otpor sredstva u kojem tijelo titra:

- ▶ Oscilator titra sve manjom amplitudom sve dok titranje potpuno ne utrne.
- ▶ Energija se gubi u okolinu.
- ▶ Najjednostavnije, otpor ovisi o brzini: $F_0 = -bv$ ($v = ds/dt$).
- ▶ Jednadžba gibanja: $m \frac{d^2s}{dt^2} + b \frac{ds}{dt} + ks = 0$
 - $2\gamma = \frac{b}{m}$ (jakost sile otpora sredstva)
 - $\omega_0 = \sqrt{\frac{k}{m}}$ (vlastita kružna frekvencija)

$$\frac{d^2s}{dt^2} + 2\gamma \frac{ds}{dt} + \omega_0^2 s = 0$$

Prigušeno titranje

$$\frac{d^2s}{dt^2} + 2\gamma \frac{ds}{dt} + \omega_0^2 s = 0 \quad \text{rj.:} \quad s = A_0 e^{-\gamma t} \sin(\omega t + \varphi_0)$$



- $A = A_0 e^{-\gamma t}$ amplituda se eksponencijalno smanjuje s vremenom.
- $\omega = \sqrt{\omega_0^2 - \gamma^2}$ frekvencija nešto manja od vlastite, $\omega_0 \rightarrow$ period T veći od T_0 .

Prisilne oscilacije

- ▶ Izgubljena energija prigušenog titranja može se nadoknaditi vanjskom periodičkom silom:

$$F = F_0 \cos \omega_f t$$

- ▶ Jednadžba gibanja: $m \frac{d^2 s}{dt^2} + b \frac{ds}{dt} + ks = F_0 \cos \omega_f t$

- ▶ Sustav titra nametnutom frekvencijom: $s = A \sin(\omega_f t + \varphi_f)$

- ▶ Amplituda je stalna - neprigušeno titranje, ali ovisi o frekvenciji vanjske sile:

$$A = A(\omega_f) = \frac{\frac{F_0}{m}}{\sqrt{(\omega_f^2 - \omega_0^2)^2 + 4\gamma^2 \omega_f^2}}$$



Prisilne oscilacije

- Rezonancija - izrazito povećanje amplitude na određenoj frekvenciji.

